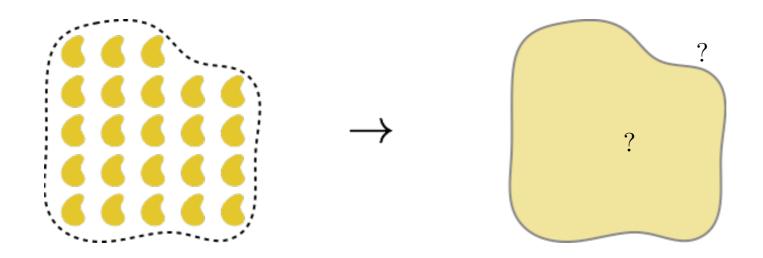
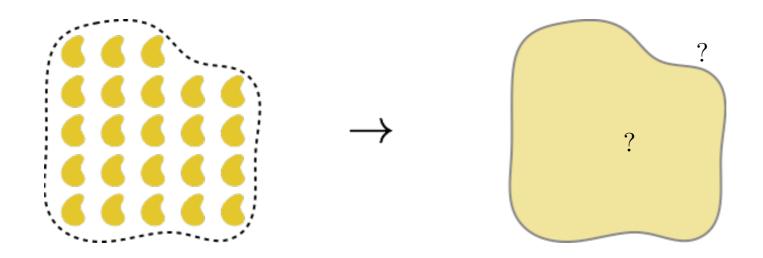


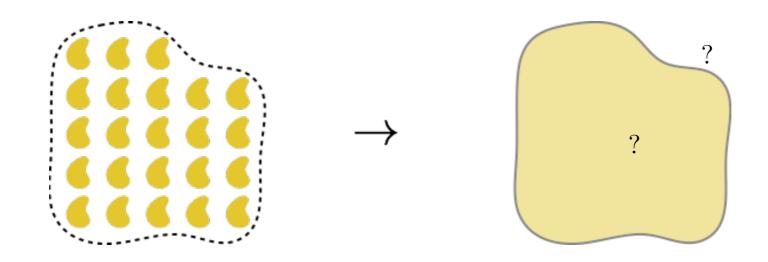
Homogeneous medium



Homogeneous medium



 ${\bf Homogeneous\ medium}$   ${\bf interface}$ 



Homogeneous medium interface

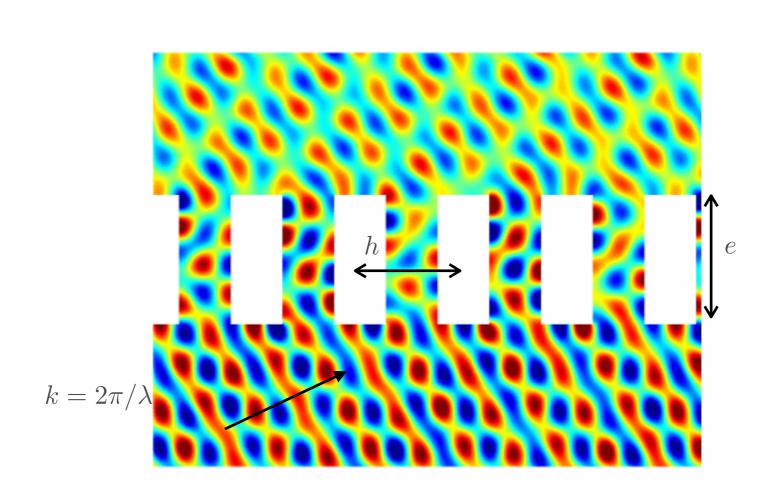
- 1) For which kind of structures?
- 2) Propagation and boundary layer effects
- 3) Examples (local resonance or not)

1) For which kind of structures?

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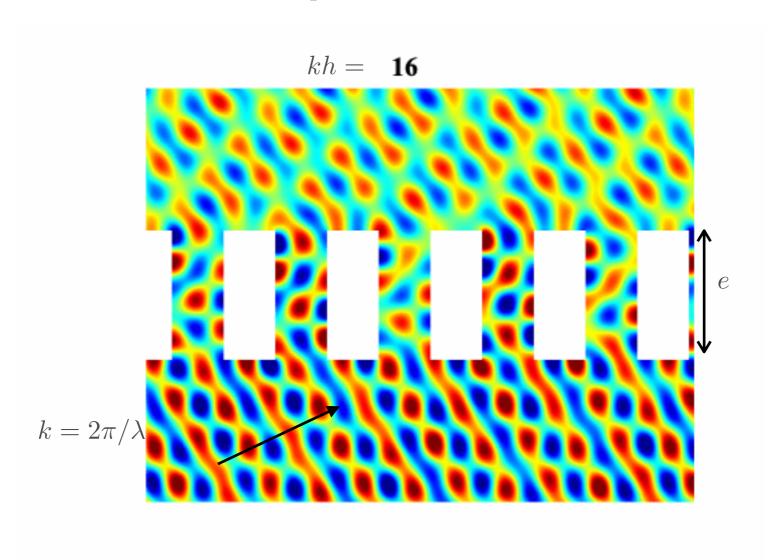
1) For which kind of structures?

periodic structures

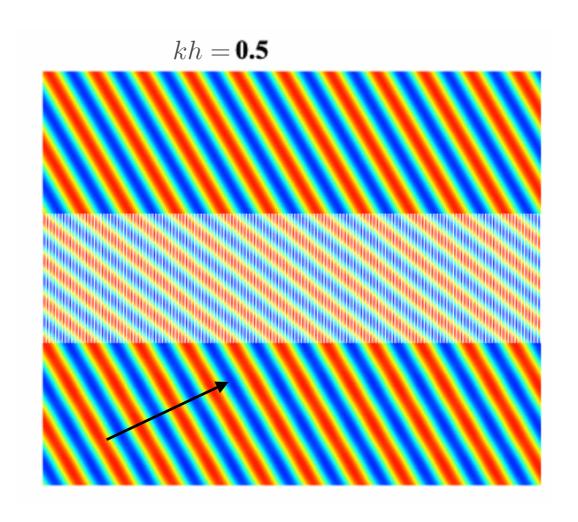


1) For which kind of structures?

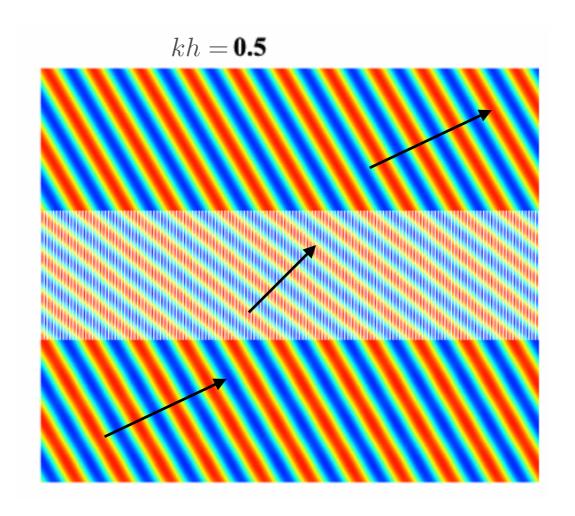
periodic structures



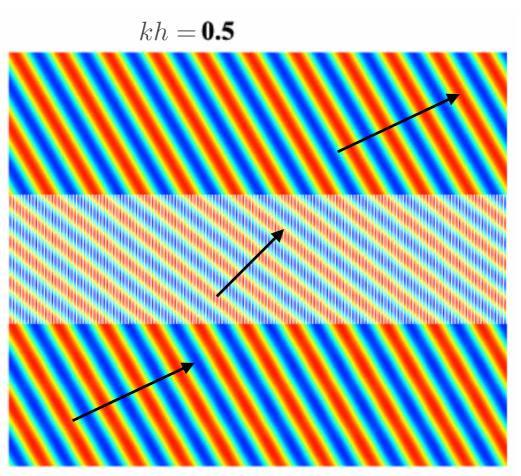
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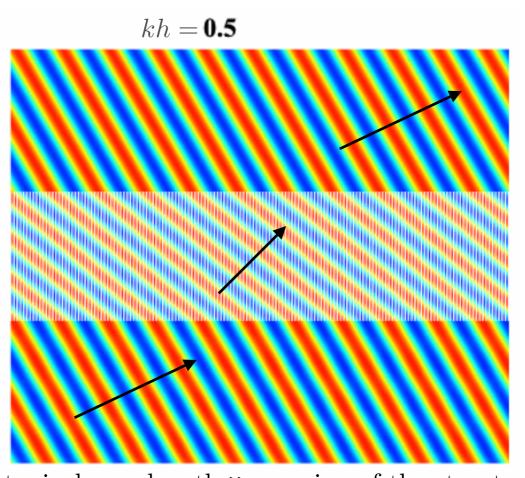
1) For which kind of structures?



typical wavelength  $\gg$  spacing of the structure  $\lambda = 2\pi/k$ 

1) For which kind of structures?

periodic structures

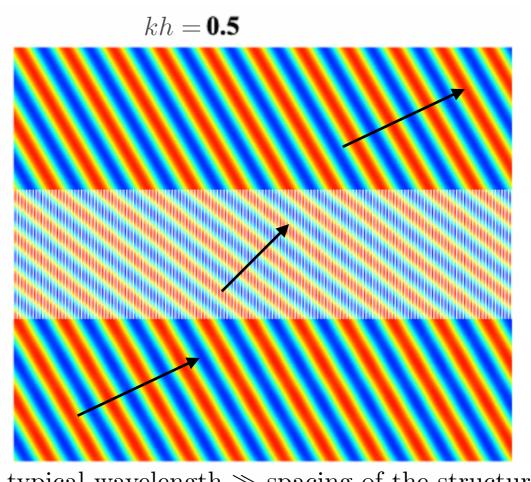


typical wavelength  $\gg$  spacing of the structure  $\lambda = 2\pi/k$ 

for  $kh \ll 1$ 

1) For which kind of structures?

periodic structures

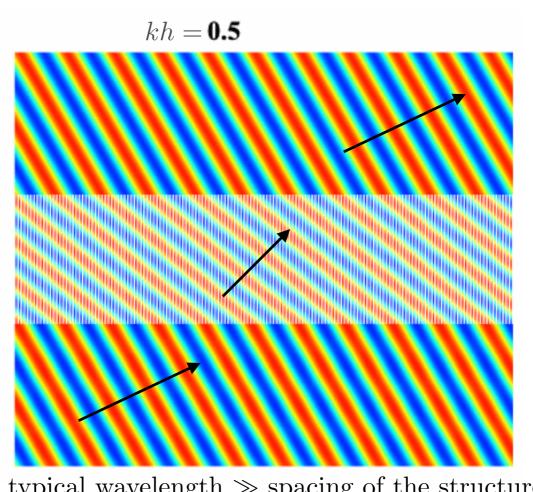


typical wavelength  $\gg$  spacing of the structure  $\lambda = 2\pi/k$ 

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1) For which kind of structures?

periodic structures



typical wavelength  $\gg$  spacing of the structure  $\lambda = 2\pi/k$ 

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Other strategies for finite wavelengths  $\to$  Course of Bojan Guzina Other strategies for quasiperiodic structures  $\to$  Course of Sébastien Guenneau

1) For which kind of structures?

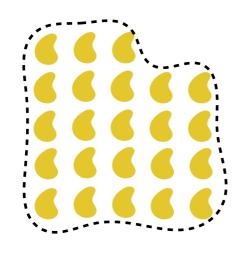
1) For which kind of structures?

2) Propagation and boundary layer effects

In its most classical form, the homogenization interrogates effective propagation

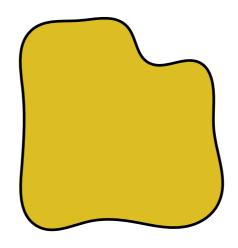
#### 2) Propagation and boundary layer effects

In its most classical form, the homogenization interrogates effective propagation



$$\operatorname{div}(a\nabla p) - b \,\frac{\partial^2 p}{\partial t^2} = 0$$

a and b vary in space

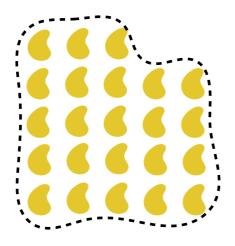


$$\operatorname{div}(\mathsf{a}_{\mathrm{eff}}\nabla p) - b_{\mathrm{eff}} \; \frac{\partial^2 p}{\partial t^2} = 0$$

effective parameters  $\mathbf{a}_{\rm eff}$  (tensor) and  $b_{\rm eff}$  (scalar) constant in space

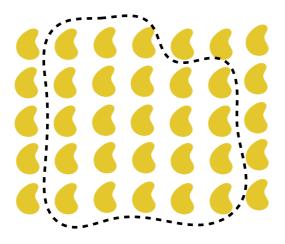
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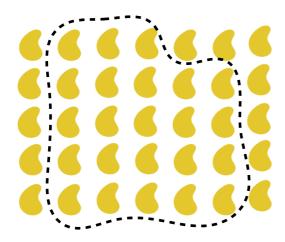
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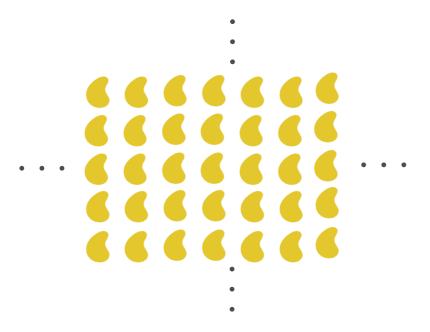
In its most classical form, the homogenization interrogates effective propagation



The (classical) homogenization in the bulk assumes an infinite medium

2) Propagation and boundary layer effects

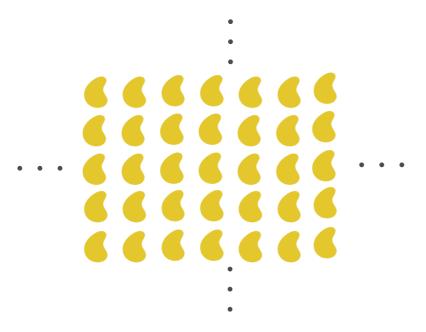
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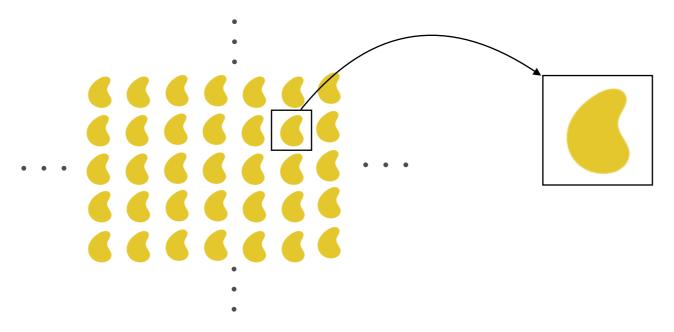
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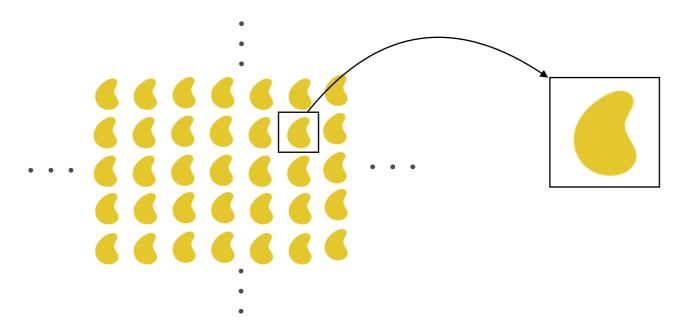
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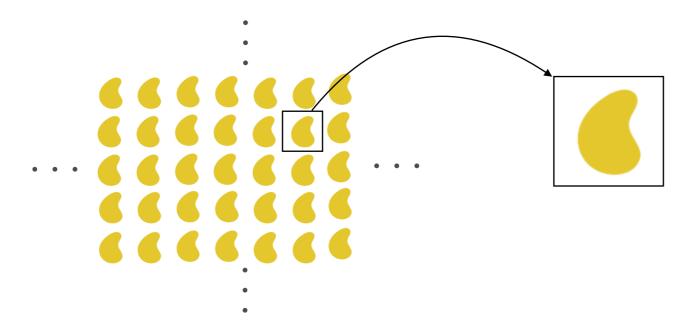
2) Propagation and boundary layer effects

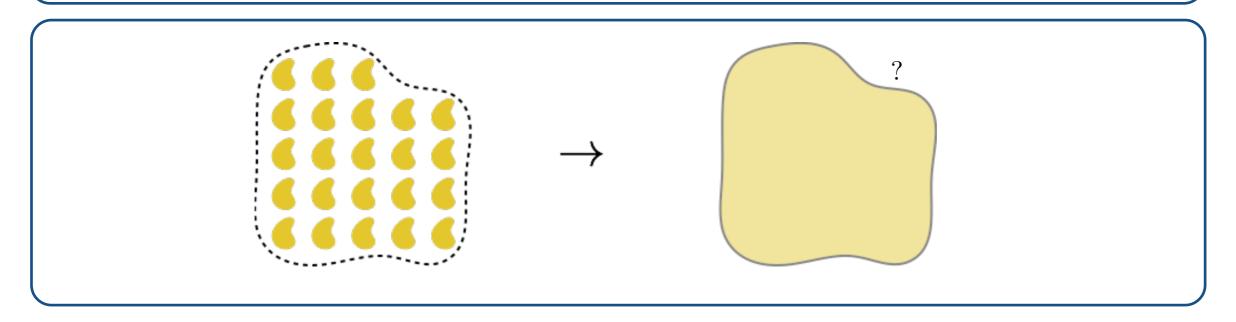
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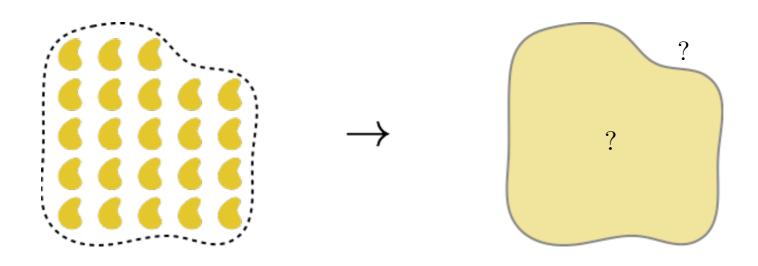


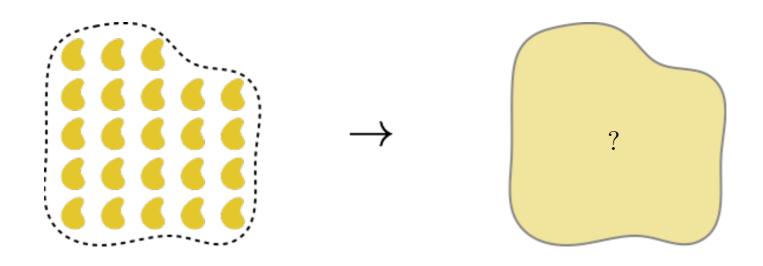
2) Propagation and boundary layer effects

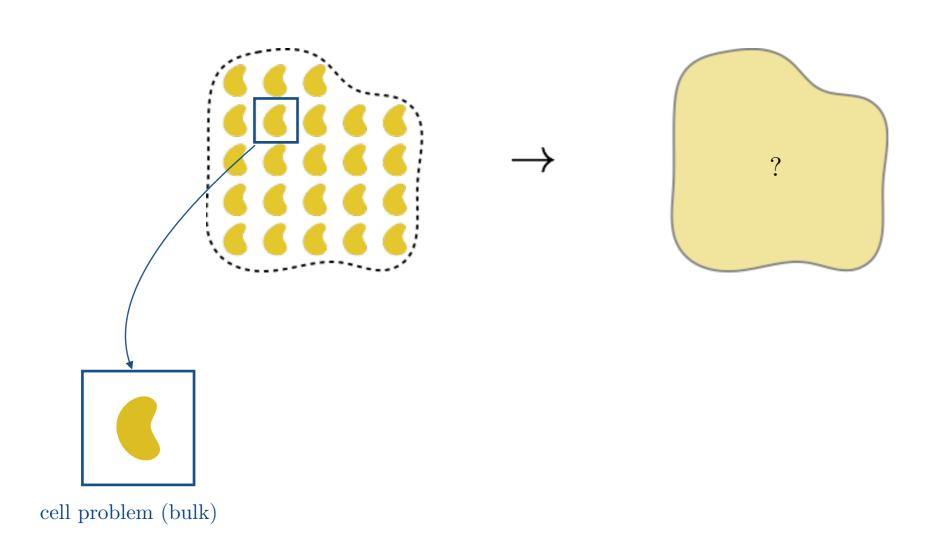
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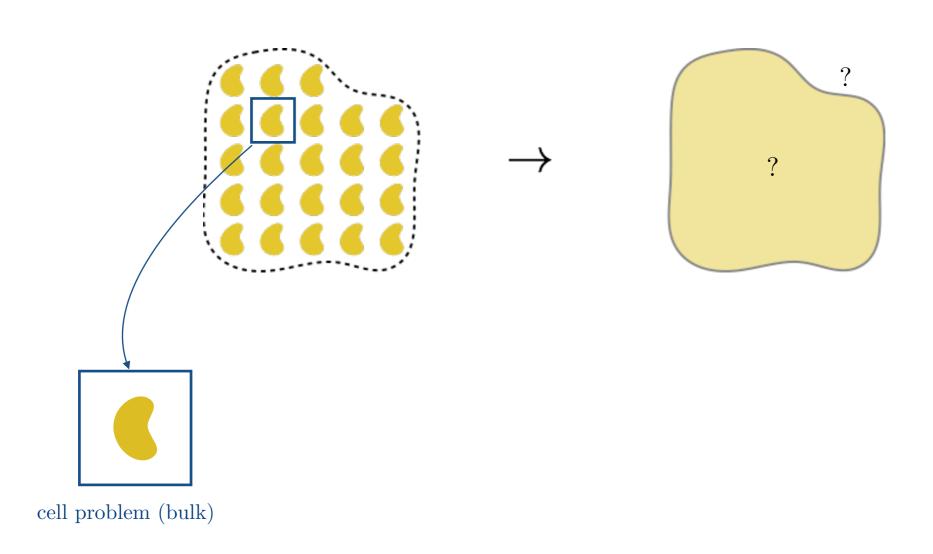


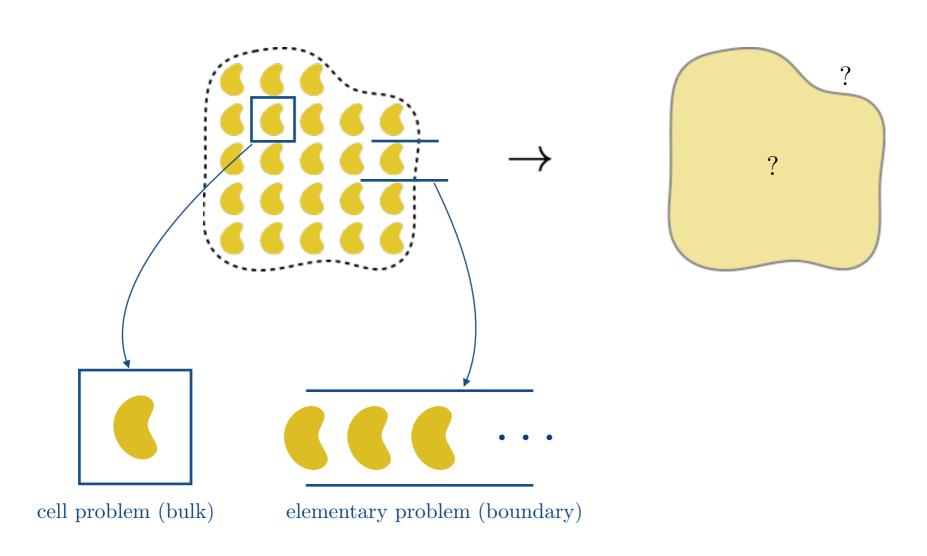




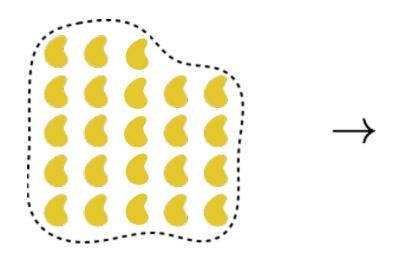






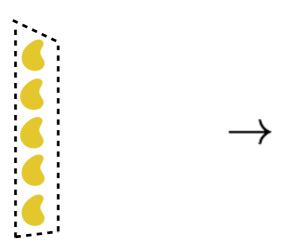


2) Propagation and boundary layer effects



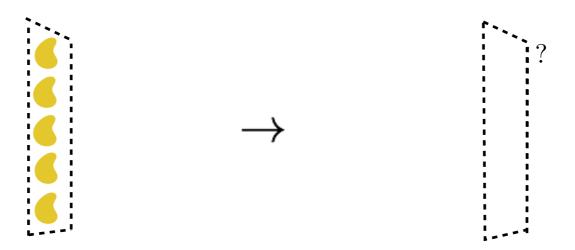
A particular case is that of small thickness of the structure

2) Propagation and boundary layer effects



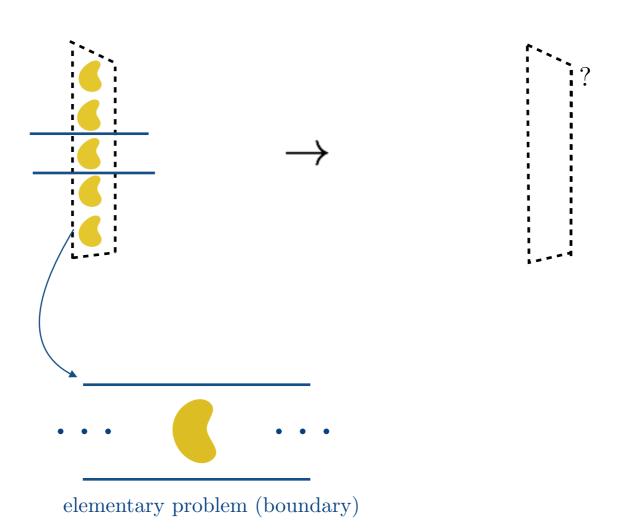
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2) Propagation and boundary layer effects



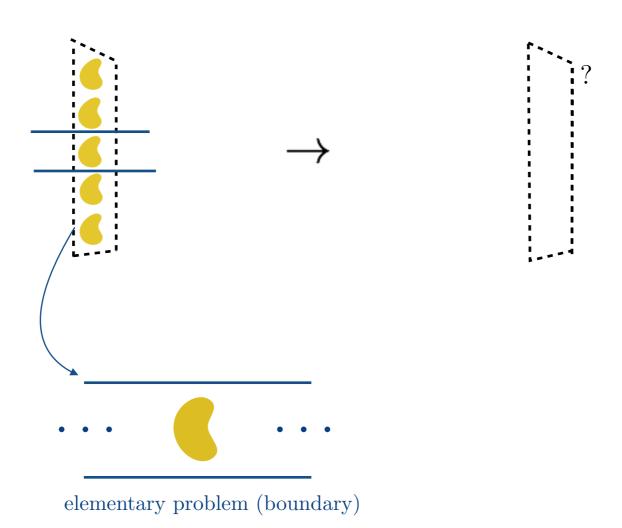
A particular case is that of small thickness of the structure

2) Propagation and boundary layer effects



A particular case is that of small thickness of the structure

2) Propagation and boundary layer effects



A particular case is that of small thickness of the structure Other strategies of homogenization  $\to$  Course of Bérangère Delourme

2) Propagation and boundary layer effects

2 different questions that have to be addressed

3) Examples (local resonance or not)

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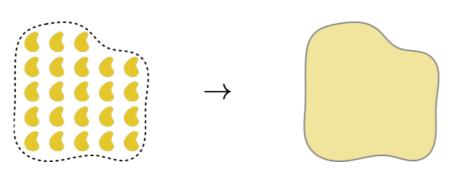
Asymptotic analysis :  $\eta = kh \rightarrow 0$ 

3) Examples (local resonance or not)

Asymptotic analysis :  $\eta = kh \rightarrow 0$ 

We can expect resonances of the resulting medium:

• if resonances take place in the resulting structure



Faraday cage, FP interferometer · · ·

 $\rightarrow$  Course of Kim Pham

 $\bullet$  if a single inclusion supports resonances



subwavelength resonance: Mie, Helmholtz, Minnaert · · ·

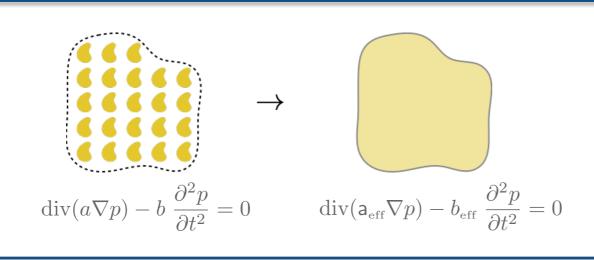
Appropriate scalings are needed to encapsulate the resonance

for instance Mie

$$kh = \eta \ll 1$$
 but  $k_0h = O(1) \rightarrow c_0/c = O(\eta)$ 

 $\rightarrow$  Course of Claude Boutin

#### Maurel & Pham



classical homogenization: effective wave equation properties of the effective parameters

Acoustic case



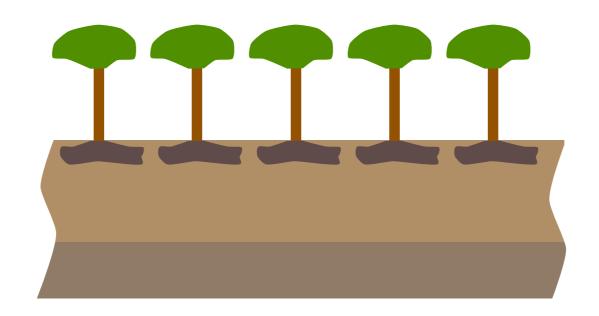
Higher order homogenization, including transmission conditions

Acoustic case



Homogenization of an array of beams on the top of an elastic half-space

2D elastic case



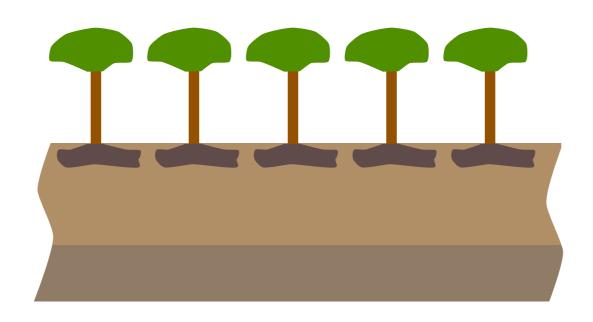
Sébastien Guenneau, Institut Fresnel, Marseille - France

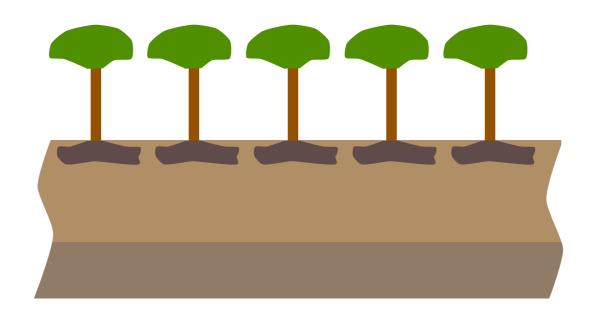
Agnès Maurel, Institut Langevin, ESPCI, Paris - France

Jean-Jacques Marigo, Laboratoire de Mécanique des Solides, Polytechnique, Palaiseau - France

Kim Pham,

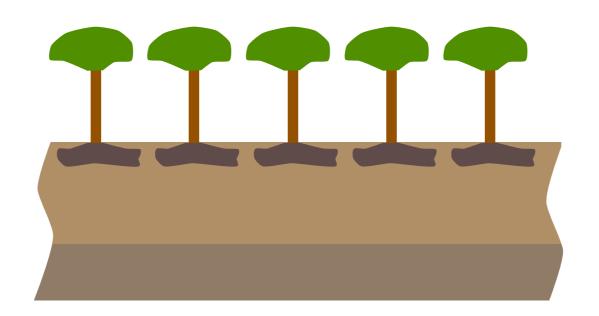
Institut des Sciences de la Mécanique et Applications Industrielles, ENSTA, Palaiseau - France





Problem setting:

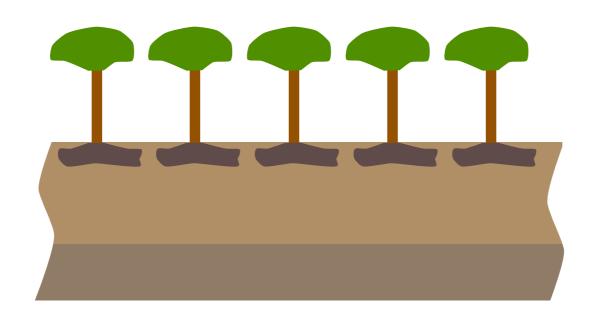
Anti-plane elasticity



Problem setting:

Anti-plane elasticity

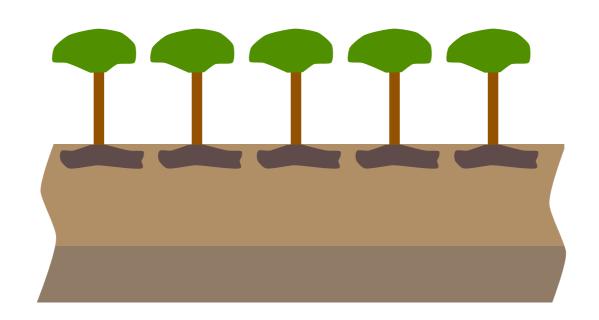
• Guiding layer over a substrate



Problem setting:

Anti-plane elasticity

- Guiding layer over a substrate
- Forest of trees

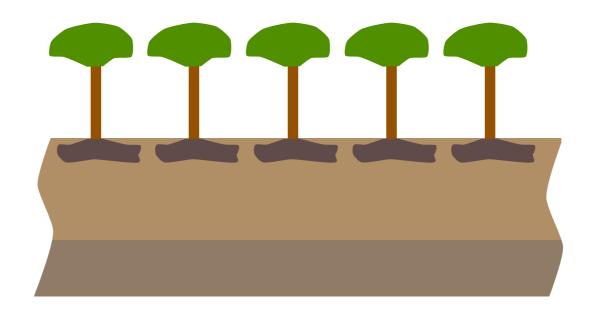


Problem setting:

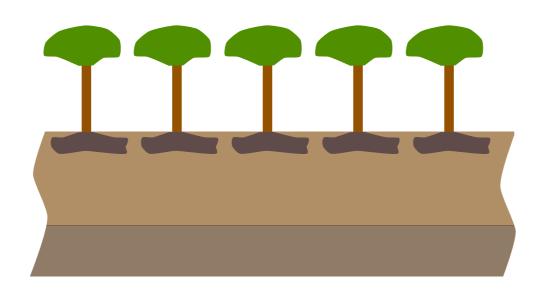
Anti-plane elasticity

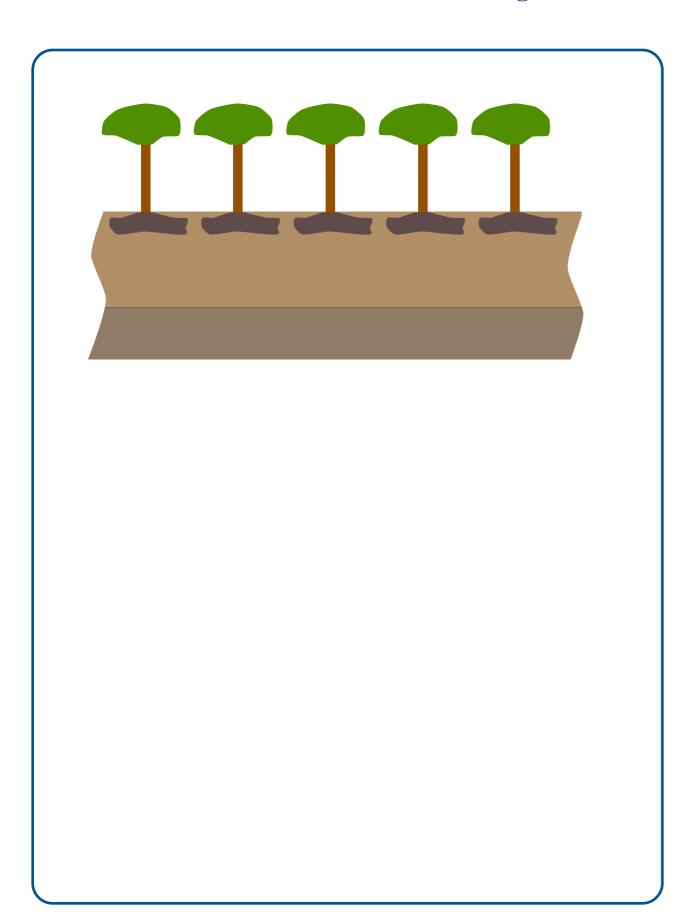
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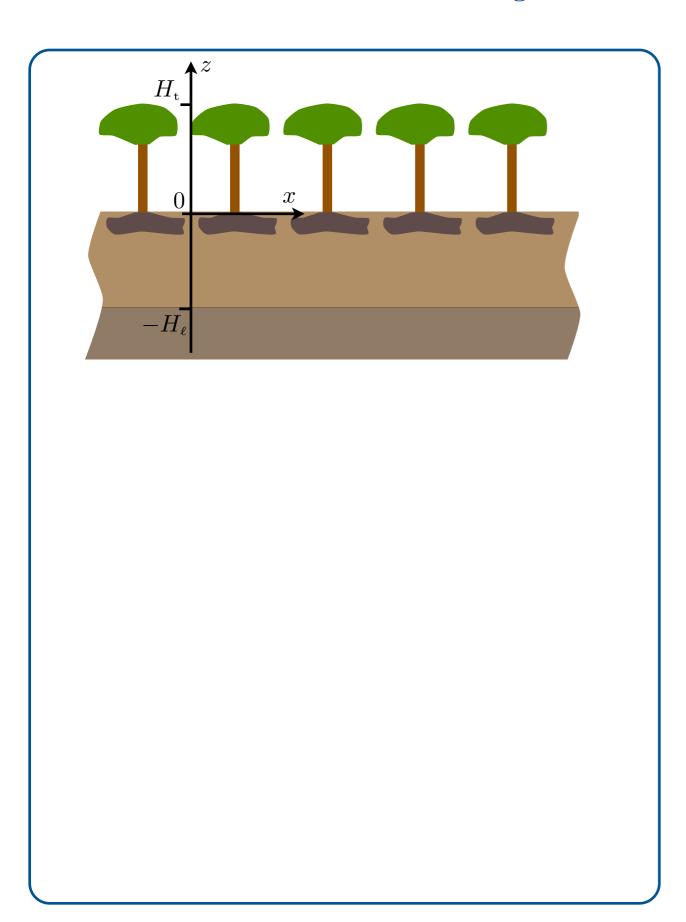
How do Love waves propagate and interact with trees?

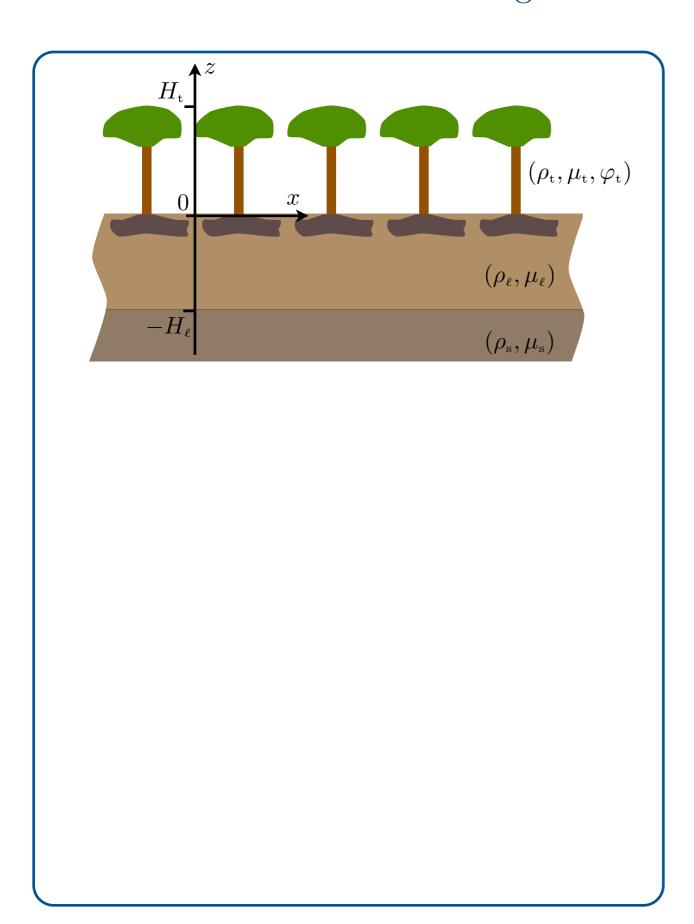


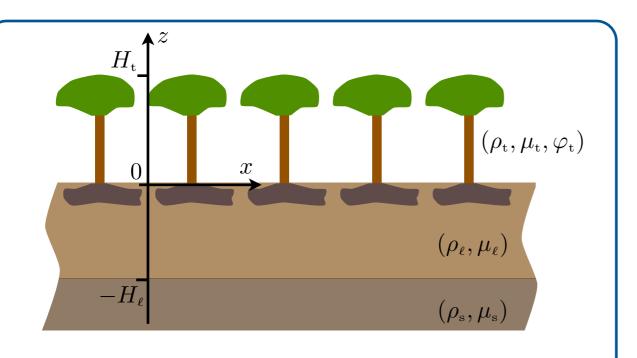
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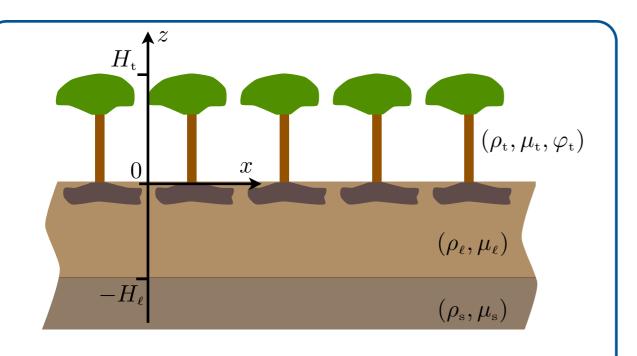






in the trees, layer, substrate

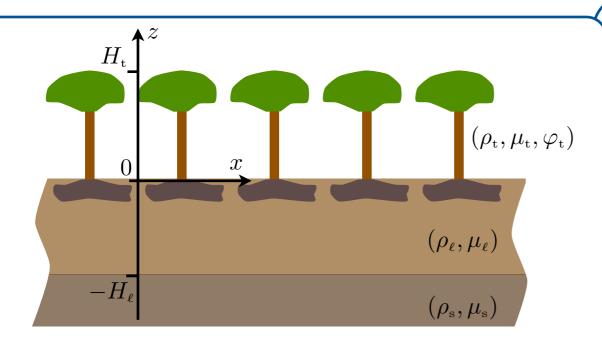
$$\sigma = \mu \nabla u, \quad \rho \frac{\partial^2 u}{\partial t^2} = \text{div} \sigma,$$



in the trees, layer, substrate

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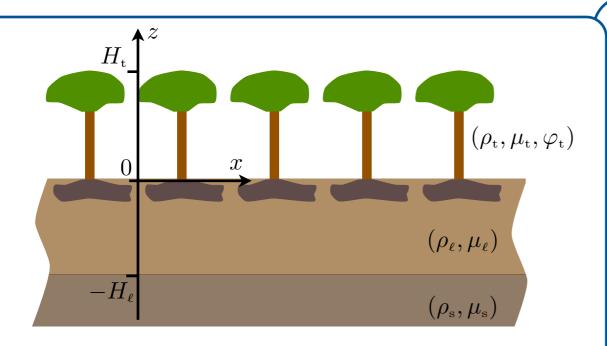
continuity of u and  $\sigma \cdot \mathbf{n}$  between two elastic media Neumann b.c.  $\sigma \cdot \mathbf{n} = 0$  elastic medium/air



in the trees, layer, substrate

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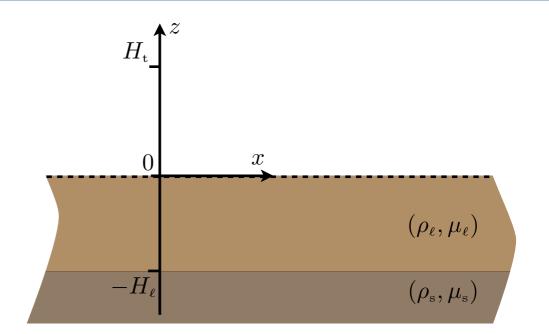
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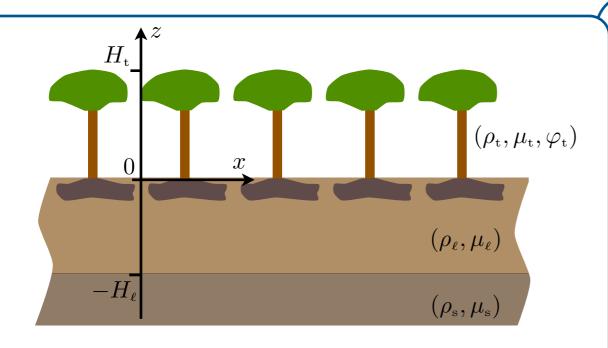
in the trees, layer, substrate

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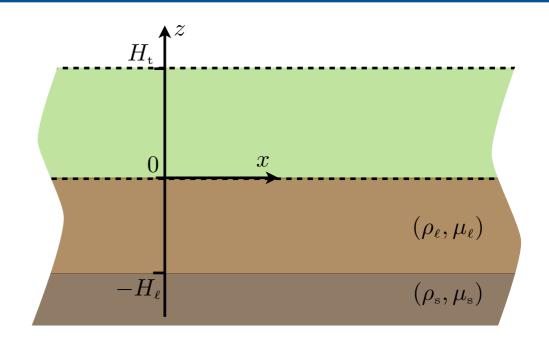
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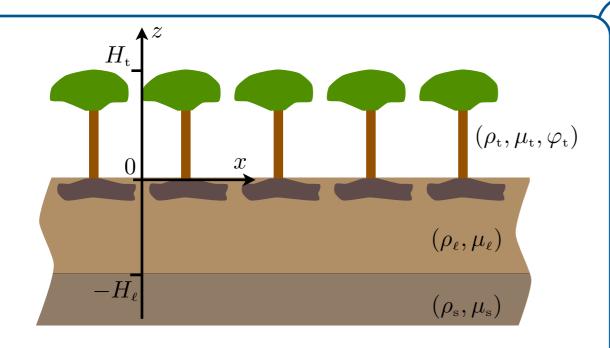
continuity of u and  $\sigma \cdot \mathbf{n}$  between two elastic media Neumann b.c.  $\sigma \cdot \mathbf{n} = 0$  elastic medium/air



in the effective medium (trees)

$$oldsymbol{\sigma} = \mu_{\mathrm{t}} \left( egin{array}{cc} 0 & 0 \ 0 & arphi_{\mathrm{t}} \end{array} 
ight) 
abla u, \quad 
ho_{\mathrm{t}} rac{\partial^2 u}{\partial t^2} = \mathrm{div} oldsymbol{\sigma},$$

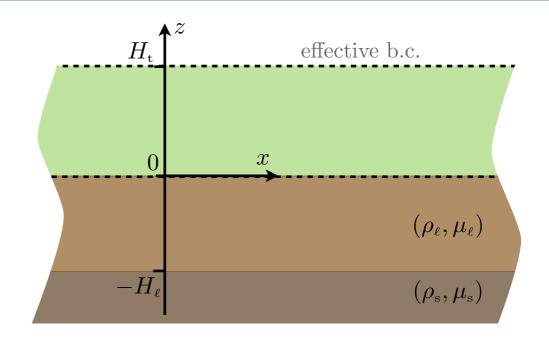
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in the trees, layer, substrate

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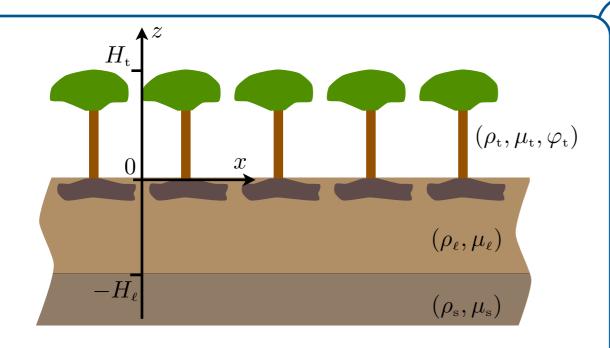
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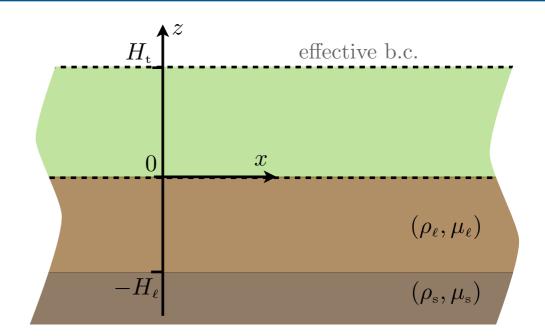
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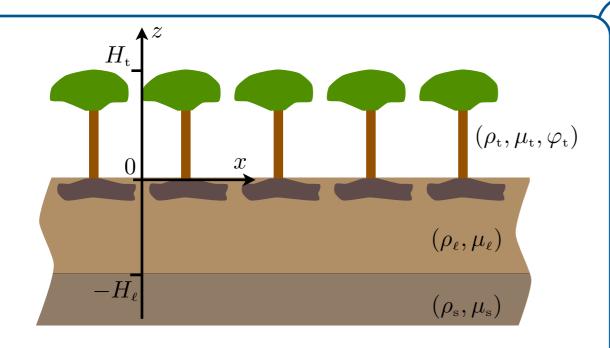


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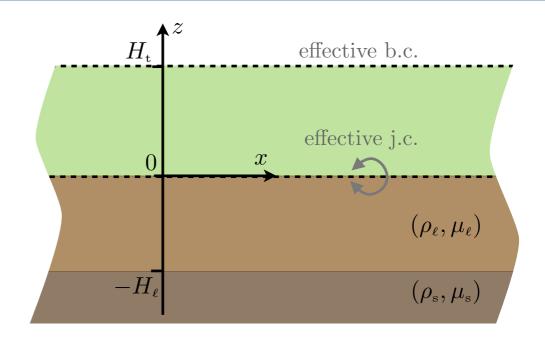
• 
$$z = H_{\rm t},$$
  $\sigma_z = -L_{\rm e} \, \frac{\partial \sigma_z}{\partial z}$ 



in the trees, layer, substrate

$$\boldsymbol{\sigma} = \mu \nabla u, \quad \rho \frac{\partial^2 u}{\partial t^2} = \operatorname{div} \boldsymbol{\sigma},$$

continuity of u and  $\boldsymbol{\sigma} \cdot \mathbf{n}$  between two elastic media Neumann b.c.  $\boldsymbol{\sigma} \cdot \mathbf{n} = 0$  elastic medium/air

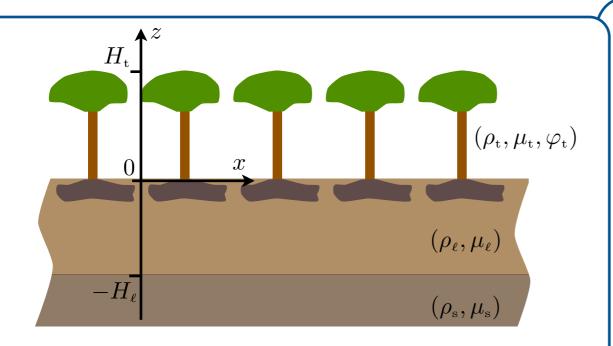


in the effective medium (trees)

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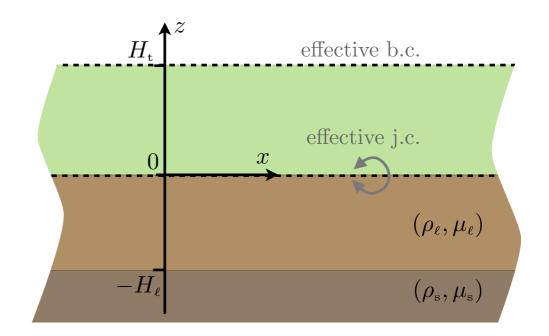
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in the trees, layer, substrate

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continuity of u and  $\boldsymbol{\sigma} \cdot \mathbf{n}$  between two elastic media Neumann b.c.  $\boldsymbol{\sigma} \cdot \mathbf{n} = 0$  elastic medium/air



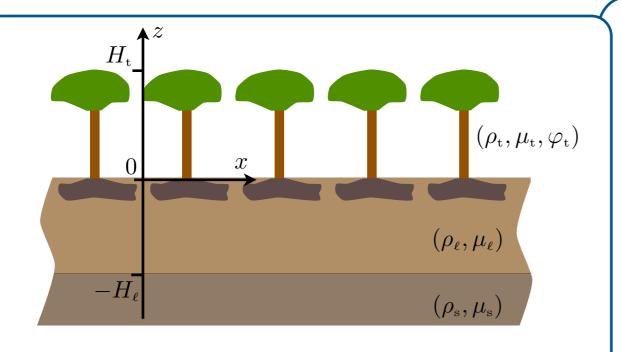
in the effective medium (trees)

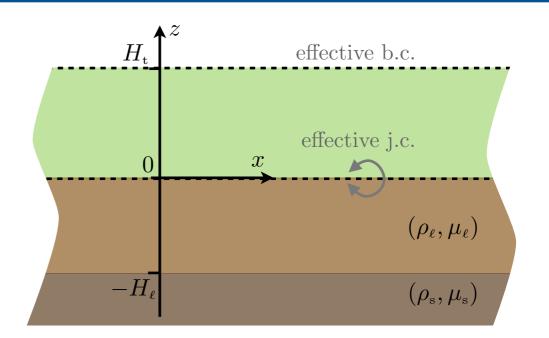
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• 
$$z = H_{\rm t},$$
  $\sigma_z = -L_{\rm e} \, \frac{\partial \sigma_z}{\partial z}$ 

• 
$$z = 0$$
, 
$$[u] = \mathcal{B}_1 \sigma_z + \mathcal{B}_2 \frac{\partial u}{\partial x}$$
$$[\sigma_z] = \mathcal{B}_2 \frac{\partial \sigma_z}{\partial x} - \mathcal{B}_3 \frac{\partial^2 u}{\partial x^2} + \rho_e \frac{\partial^2 u}{\partial t^2}$$





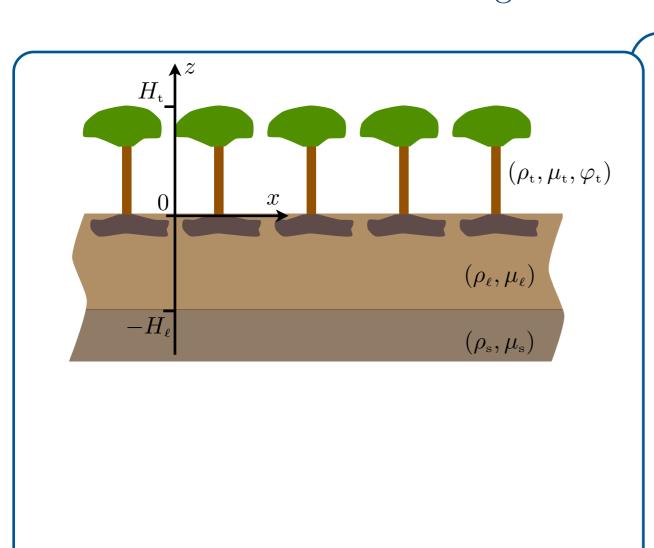
in the effective medium (trees)

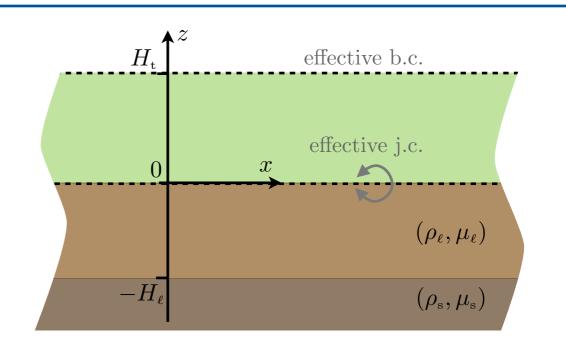
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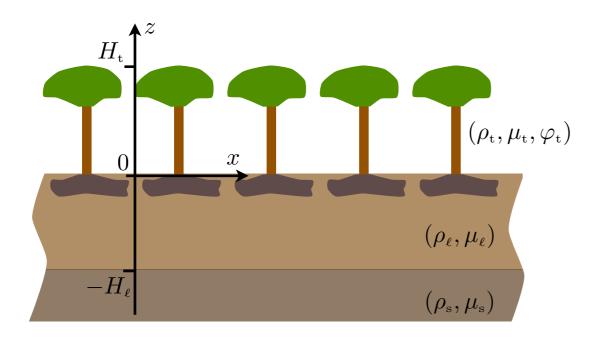
$$\boldsymbol{\sigma} = \mu \nabla u, \quad \rho \frac{\partial^2 u}{\partial t^2} = \operatorname{div} \boldsymbol{\sigma},$$

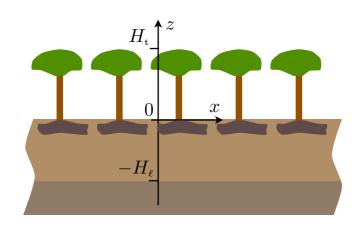
• 
$$z = H_{\rm t},$$
  $\sigma_z = -L_{\rm e} \, \frac{\partial \sigma_z}{\partial z}$ 

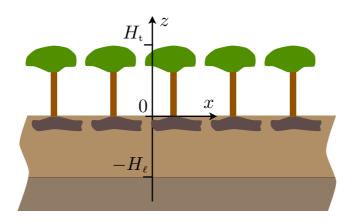
• 
$$z = 0$$
, 
$$[u] = \mathcal{B}_1 \sigma_z + \mathcal{B}_2 \frac{\partial u}{\partial x}$$
$$[\sigma_z] = \mathcal{B}_2 \frac{\partial \sigma_z}{\partial x} - \mathcal{B}_3 \frac{\partial^2 u}{\partial x^2} + \rho_e \frac{\partial^2 u}{\partial t^2}$$



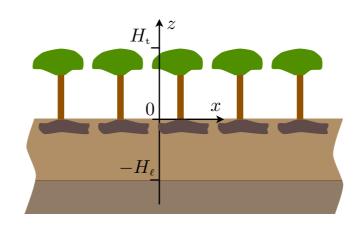




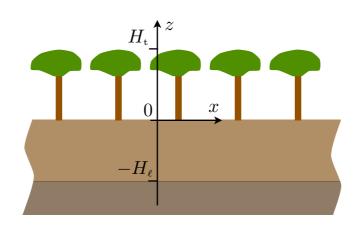




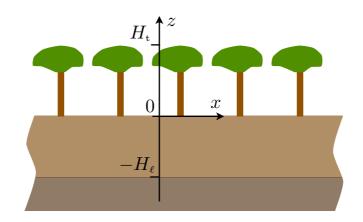
Results in the harmonic regime



Results in the harmonic regime

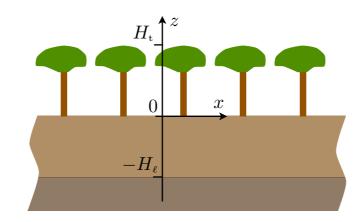


Results in the harmonic regime



Guided waves for the homogenized problem

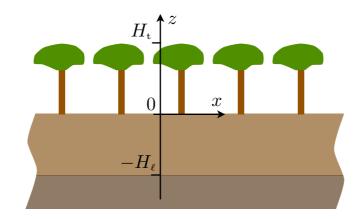
Results in the harmonic regime



Guided waves for the homogenized problem

$$u(x,z) = \begin{cases} e^{\alpha_{s}z}e^{i\beta x}, & z < -H_{\ell}, \\ (A\cos k_{\ell}z + B\sin k_{\ell}z)e^{i\beta x}, & -H_{\ell} < z < 0 \\ (C\cos k_{t}z + D\sin k_{t}z)e^{i\beta x}, & 0 < z < H_{t} \end{cases}$$

Results in the harmonic regime



#### Guided waves for the homogenized problem

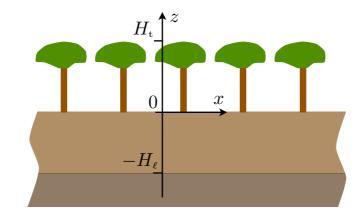
$$u(x,z) = \begin{cases} e^{\alpha_{\rm s} z} e^{i\beta x}, & z < -H_{\ell}, \\ (A\cos k_{\ell} z + B\sin k_{\ell} z) e^{i\beta x}, & -H_{\ell} < z < 0 \\ (C\cos k_{\rm t} z + D\sin k_{\rm t} z) e^{i\beta x}, & 0 < z < H_{\rm t} \end{cases}$$

$$\alpha_{s} = \sqrt{\beta^{2} - \frac{\omega^{2}}{c_{s}^{2}}}$$

$$k_{\ell} = \sqrt{\frac{\omega^{2}}{c_{\ell}^{2}} - \beta^{2}}$$

$$k_{t} = \frac{\omega}{c_{t}}$$

Results in the harmonic regime



Guided waves for the homogenized problem

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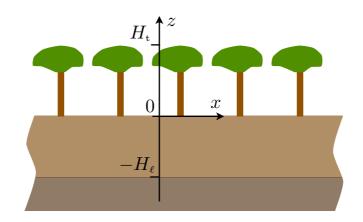
$$k_{\ell} = \sqrt{\frac{\omega^{2}}{c_{\ell}^{2}} - \beta^{2}}$$

$$k_{t} = \frac{\omega}{c_{t}}$$

Dispersion relation:

$$1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{s} \alpha_{s}} \tan k_{\ell} H_{\ell} - \frac{\mu_{t} k_{t}}{\mu_{\ell} k_{\ell}} \varphi_{t} \tan k_{t} H_{t} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{s} \alpha_{s}} \right) = 0.$$
 (without foliage)

Results in the harmonic regime



$$\alpha_{\rm s} = \sqrt{\beta^2 - \frac{\omega^2}{c_{\rm s}^2}}$$

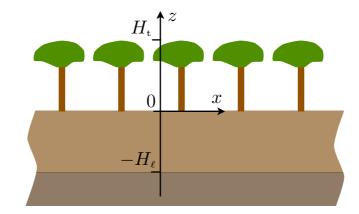
$$k_{\ell} = \sqrt{\frac{\omega^2}{c_{\ell}^2} - \beta^2}$$

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 (without foliage)

Results in the harmonic regime



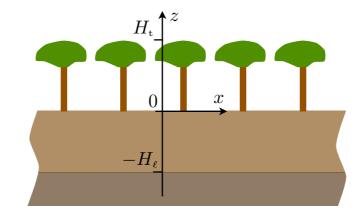
$$\alpha_{\rm s} = \sqrt{\beta^2 - \frac{\omega^2}{c_{\rm s}^2}}$$

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$$1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{s} \alpha_{s}} \tan k_{\ell} H_{\ell} - \frac{\mu_{t} k_{t}}{\mu_{\ell} k_{\ell}} \varphi_{t} \tan k_{t} H_{t} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{s} \alpha_{s}} \right) = 0.$$
 (without foliage)

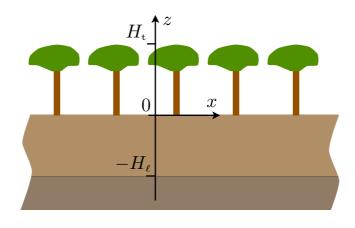


$$\text{Dispersion relation:} \qquad 1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \tan k_{\ell} H_{\ell} - \frac{\mu_{\mathrm{t}} k_{\mathrm{t}}}{\mu_{\ell} k_{\ell}} \varphi_{\mathrm{t}} \tan k_{\mathrm{t}} H_{\mathrm{t}} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \right) = 0.$$

$$\alpha_{\rm s} = \sqrt{\beta^2 - \frac{\omega^2}{c_{\rm s}^2}}$$

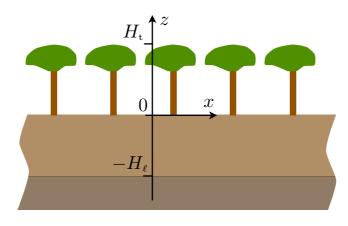
$$k_{\ell} = \sqrt{\frac{\omega^2}{c_{\ell}^2} - \beta^2}$$

$$k_{\rm t} = \frac{\omega}{c_{\rm t}}$$



$$\text{Dispersion relation:} \qquad 1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \tan k_{\ell} H_{\ell} - \frac{\mu_{\mathrm{t}} k_{\mathrm{t}}}{\mu_{\ell} k_{\ell}} \varphi_{\mathrm{t}} \tan k_{\mathrm{t}} H_{\mathrm{t}} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \right) = 0.$$

$$lpha_{
m s} = \sqrt{eta^2 - rac{\omega}{c_{
m s}^2}}$$
 $k_{
m \ell} = \sqrt{rac{\omega^2}{c_{
m \ell}^2} - eta}$ 
 $k_{
m t} = rac{\omega}{c_{
m \ell}}$ 

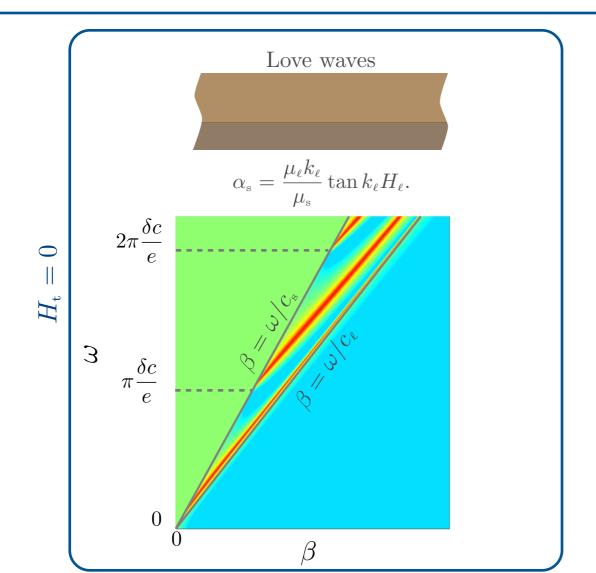


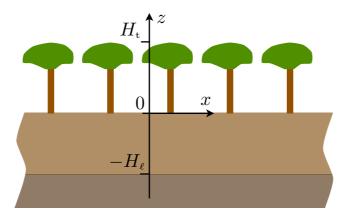
$$\text{Dispersion relation:} \qquad 1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \tan k_{\ell} H_{\ell} - \frac{\mu_{\mathrm{t}} k_{\mathrm{t}}}{\mu_{\ell} k_{\ell}} \varphi_{\mathrm{t}} \tan k_{\mathrm{t}} H_{\mathrm{t}} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \right) = 0.$$

$$\alpha_{\rm s} = \sqrt{\beta^2 - \frac{\omega^2}{c_{\rm s}^2}}$$

$$k_{\ell} = \sqrt{\frac{\omega^2}{c_{\ell}^2} - \beta}$$

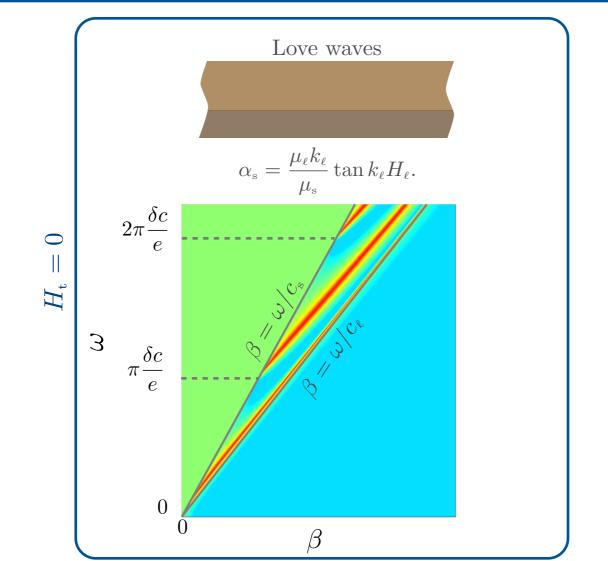
$$k_{\rm t} = \frac{\omega}{c_{\rm t}}$$

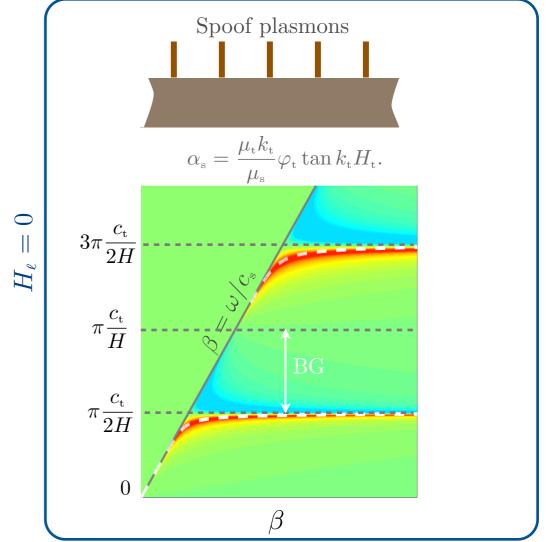


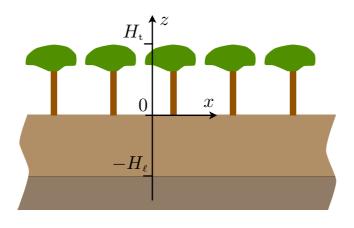


$$\text{Dispersion relation:} \qquad 1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \tan k_{\ell} H_{\ell} - \frac{\mu_{\mathrm{t}} k_{\mathrm{t}}}{\mu_{\ell} k_{\ell}} \varphi_{\mathrm{t}} \tan k_{\mathrm{t}} H_{\mathrm{t}} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \right) = 0.$$

$$lpha_{ ext{s}} = \sqrt{eta^2 - rac{\omega^2}{c_{ ext{s}}^2}}$$
 $k_{ ext{\ell}} = \sqrt{rac{\omega^2}{c_{ ext{\ell}}^2} - eta^2}$ 
 $k_{ ext{t}} = rac{\omega}{c_{ ext{t}}}$ 

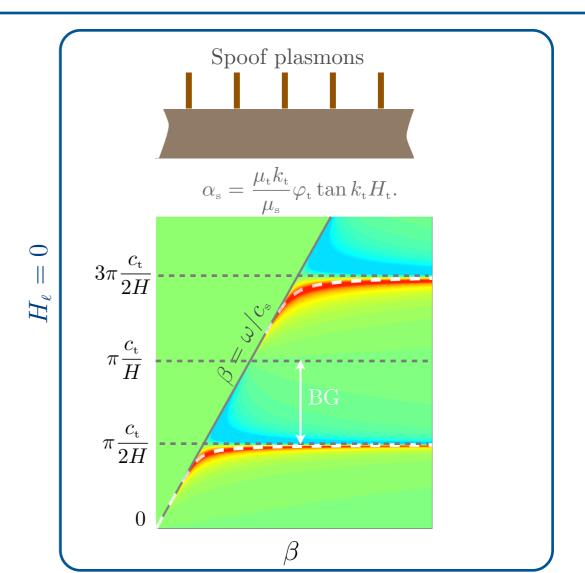


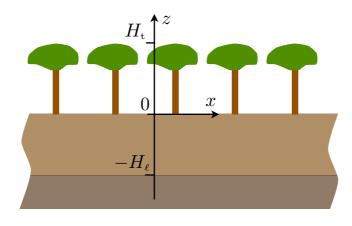




$$k_{\ell}$$
 :  $k_{
m t}$ 

$$\text{Dispersion relation:} \qquad 1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \tan k_{\ell} H_{\ell} - \frac{\mu_{\mathrm{t}} k_{\mathrm{t}}}{\mu_{\ell} k_{\ell}} \varphi_{\mathrm{t}} \tan k_{\mathrm{t}} H_{\mathrm{t}} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \right) = 0.$$



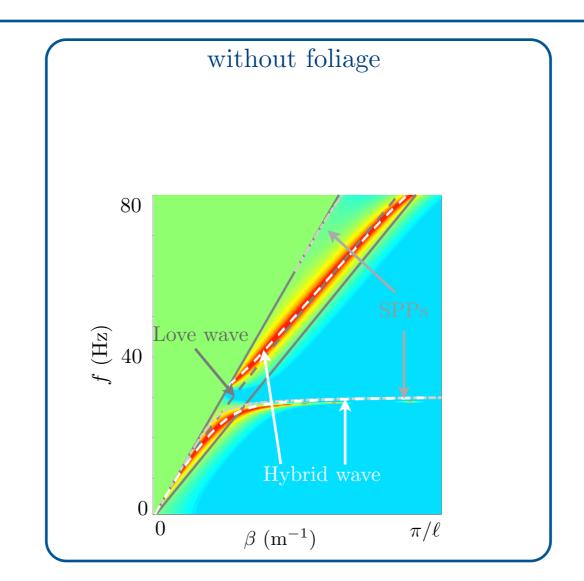


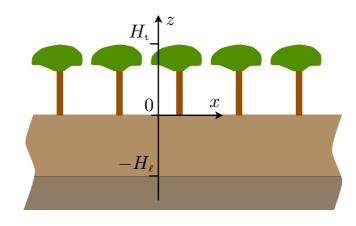
$$\text{Dispersion relation:} \qquad 1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \tan k_{\ell} H_{\ell} - \frac{\mu_{\mathrm{t}} k_{\mathrm{t}}}{\mu_{\ell} k_{\ell}} \varphi_{\mathrm{t}} \tan k_{\mathrm{t}} H_{\mathrm{t}} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \right) = 0.$$

$$\alpha_{\rm s} = \sqrt{\beta^2 - \frac{\omega^2}{c_{\rm s}^2}}$$

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$$k_{\rm t} = \frac{\omega}{c_{\rm t}}$$



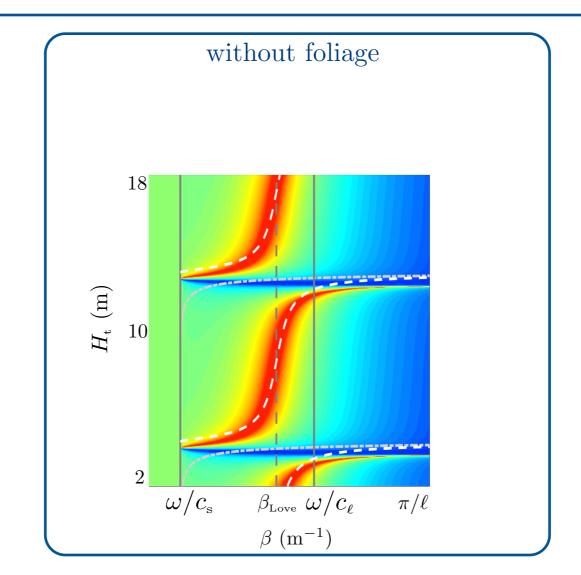


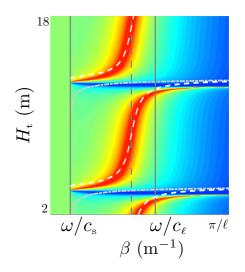
$$\text{Dispersion relation:} \qquad 1 - \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \tan k_{\ell} H_{\ell} - \frac{\mu_{\mathrm{t}} k_{\mathrm{t}}}{\mu_{\ell} k_{\ell}} \varphi_{\mathrm{t}} \tan k_{\mathrm{t}} H_{\mathrm{t}} \left( \tan k_{\ell} H_{\ell} + \frac{\mu_{\ell} k_{\ell}}{\mu_{\mathrm{s}} \alpha_{\mathrm{s}}} \right) = 0.$$

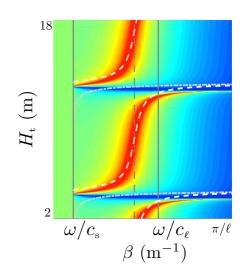
$$\alpha_{\rm s} = \sqrt{\beta^2 - \frac{\omega^2}{c_{\rm s}^2}}$$

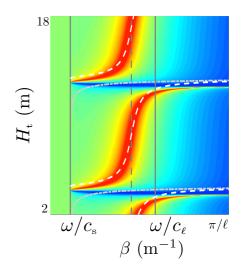
$$k_{\ell} = \sqrt{\frac{\omega^2}{c_{\ell}^2} - \beta^2}$$

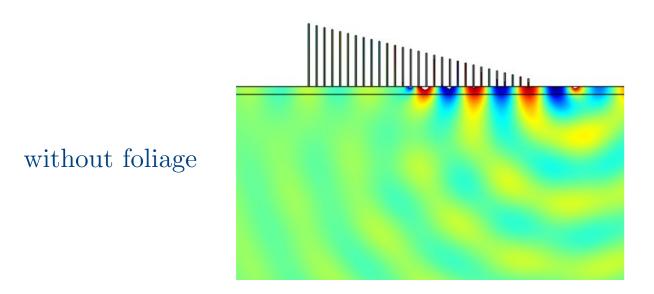
$$k_{\rm t} = \frac{\omega}{c_{\rm t}}$$

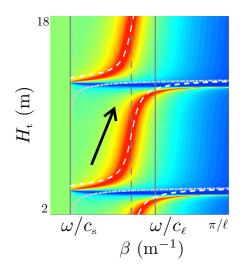


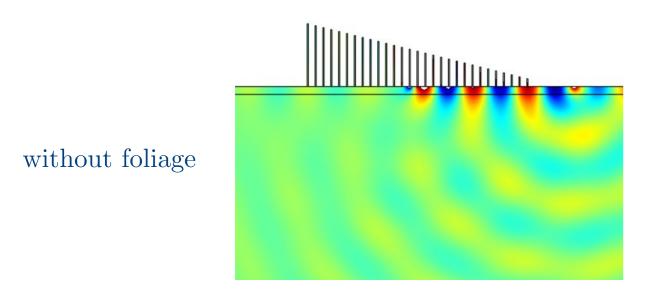


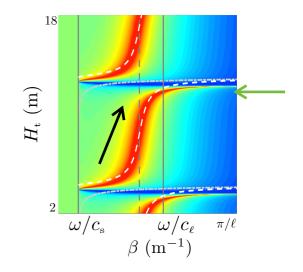


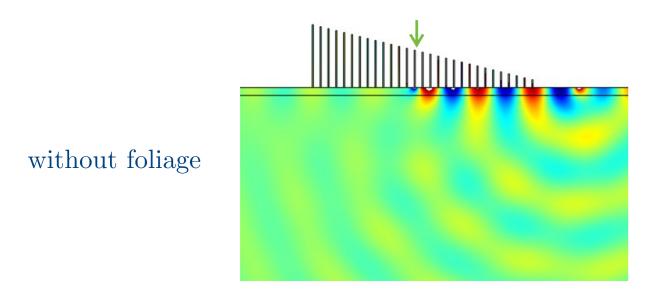


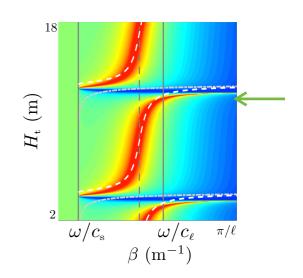


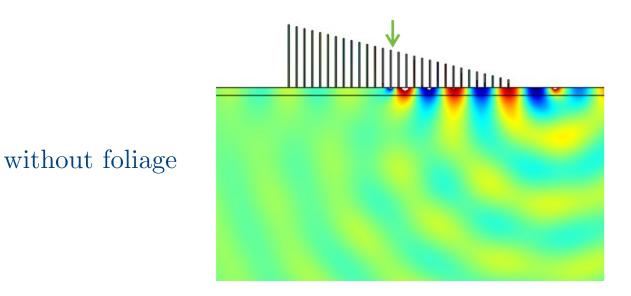


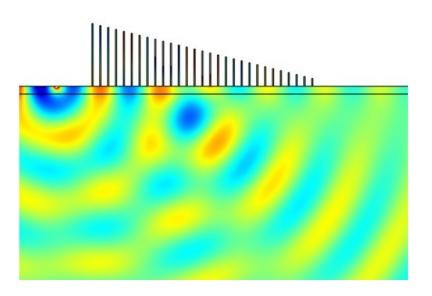


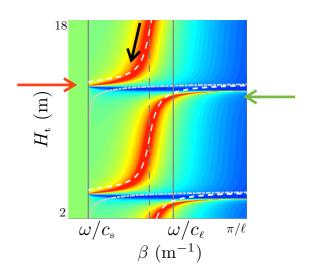


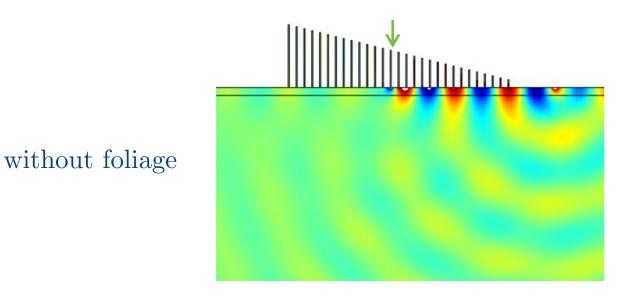


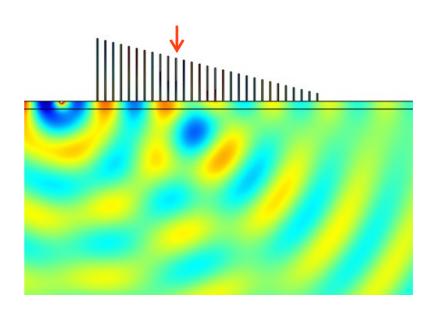


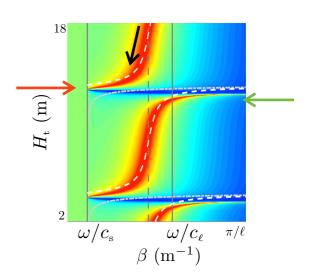


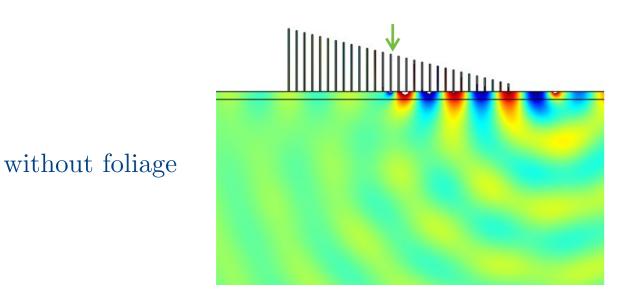


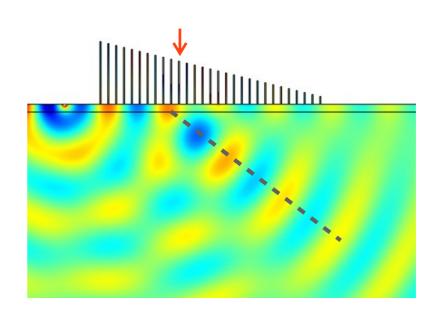


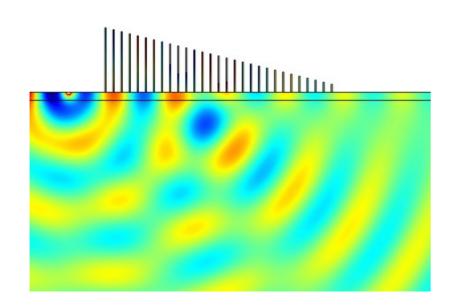


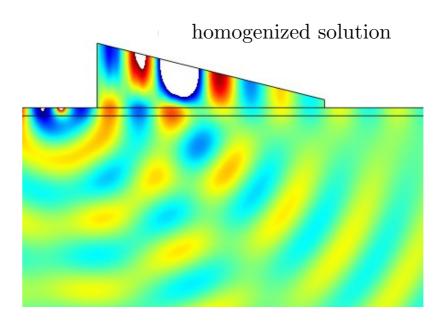






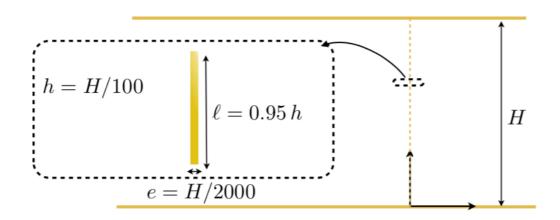


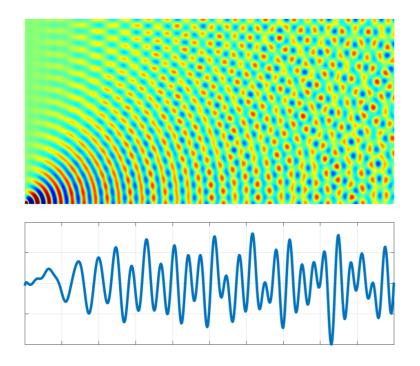




#### Asymptotic homogenization

1) For which kind of structures?





$$kH = 100, kh = 1, ke = 0.05$$
  
 $H/e = 2000$ 

